

**Erskine and Suen Reply:** While the Comment on our paper [1] by Zhong, Dong, and Xing [2] and reference to their prior related work [3] are clearly relevant to its contents and conclusions, we maintain that our experiments are not compatible with the theoretical and experimental results compiled in our Table I of Ref. [1]. We also disagree with the assertion that “better agreement” could be achieved if our data were “more properly treated.”

Zhong *et al.*'s comment generalizes a prior treatment [3] of the two dimensional Ising model to include disorder. The primary result of both models and similar models cited in our paper [4,5] that pertains to our reported experimental results is that the hysteresis loop scaling is governed by

$$A \sim A_0 + a \left( \frac{dH}{dt} \right)^\alpha, \quad (1)$$

with a (temperature dependent) constant  $A_0$  and power law exponent  $\alpha = 0.36 \pm 0.08$ . The analysis of our experimental data was based on the assumption that the scaling is governed by

$$A \sim H_0^\alpha \Omega^\beta. \quad (2)$$

Equation (1) implies  $\alpha = \beta$  in Eq. (2), and, that at very low frequencies, ( $\frac{dH}{dt} \rightarrow 0$ ), an “adiabatic” region [4] exists in which the hysteresis loss remains constant and characterized by the  $A_0$  term. Our experimental data are not compatible with  $\alpha = \beta$ . However, if we assume scaling of the form Eq. (1) with  $A_0 \neq 0$ , and fit our data allowing  $A_0$ ,  $a$ , and  $\alpha$  to be unrestricted parameters, we find values of  $\alpha$  for Eq. (1) that are essentially compatible with the corresponding parameter  $\beta$  in our paper ( $\beta \sim 0.1$ ) and not  $\beta \sim 0.35$ . A typical fit for the 3 ML thick stepped surface data at  $T = 295$  and 98 K (Fig. 3 of Ref. [1]) along with the parameters are displayed in our accompanying Fig. 1. Clearly  $A_0$  is finite and decreases with increasing temperature, as the Ising model calculations predict. However, the value of  $\alpha$  in Eq. (1) remains small and consistent with the value of  $\beta$  obtained using the scaling of Eq. (2). We again conclude that  $\alpha \neq \beta$  and that our new results for  $p(1 \times 1)$  Fe on W(110) are not compatible with results compiled in Table I of Ref. [1]. New experimental results for Co on Cu(100) and additional discussions of hysteresis scaling in ferromagnetic thin films will be presented in a forthcoming publication [6].

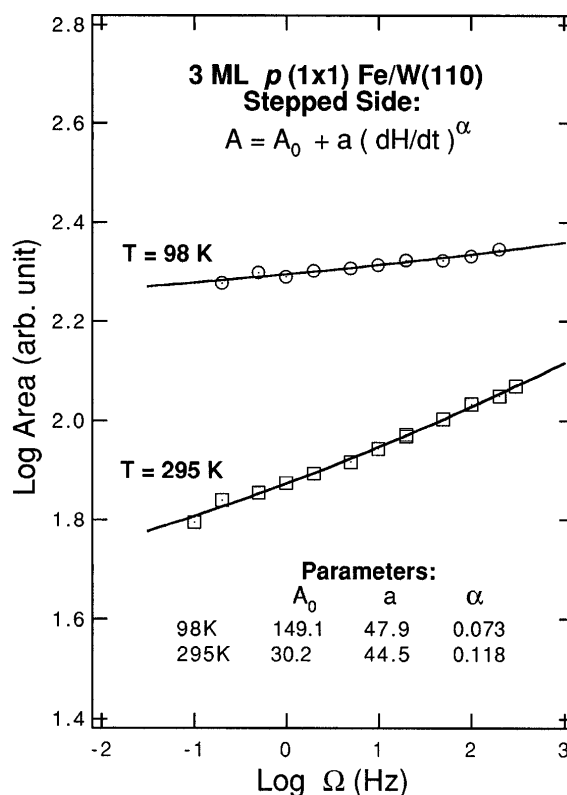


FIG. 1. Fit of experimental data of Ref. [1] to alternate scaling law [Eq. (1)]. The power law scaling exponents  $\alpha$  are larger than corresponding exponents  $\beta$  obtained from the scaling law [Eq. (2)], i.e.,  $\beta = 0.076 \pm 0.002$  for  $T = 295$  K, but are still not near the calculated value  $\alpha = 0.36 \pm 0.08$ .

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