SummaryL of unit 2 (Ch17, Ch18, Ch22.6, update: 10/14/12)

Constant: $\frac{\mu_0}{4\pi} = 10^{-7}S.I.$ Integrals: $\int_{\alpha_1}^{\alpha_2} sin\alpha d\alpha = (cos\alpha_1 - cos\alpha_2), \int_a^b r^n dr = \frac{r^{n+1}}{n+1} \Big|_a^b, \int_a^b \frac{dr}{r} = ln(b/a).$ Electric Potential: For 1 charged particle: $\Delta K + \Delta U = 0, \Delta U = -q\mathbf{E} \cdot \Delta \mathbf{l}$ (work against $F = q\mathbf{E}$). Potential: $\Delta V = \frac{\Delta U}{q} = V_f - V_i = -\mathbf{E} \cdot \Delta \mathbf{l} = -E\Delta l \cos\theta = -(E_x\Delta x + E_y\Delta y + E_z\Delta z).$ Units: V=Nm/C. From V to \mathbf{E} : $\mathbf{E} = -(\partial V/\partial x, \partial V/\partial y, \partial V/\partial z) \equiv -\nabla V.$ For spherically symm. case: $E_r = -dV(r)/dr.$ Potential diff: $V_f - V_i = -\int_i^f \mathbf{E} \cdot d\mathbf{l}.$ For Q at origin, define $V(\infty) = 0.$ $V(r) = -\int_{\infty}^r \frac{kQ}{r'^2} dr' = kQ/r.$ For a potential function, path integral along any closed loop $\Delta V = 0$. This implies conservation of energy. Potential energy U_A : work required to move q from ∞ to A, assume $U_{\infty} = 0$. Potential: $V_A = \frac{U_A}{q}, V_{\infty} = 0.$ System of many charged particles: $U = \Sigma_{i < j} U_{ij}$, where $U_{ij} = k\frac{q_i q_j}{r_{ij}}, r_{ij} = |\mathbf{r}_i - \mathbf{r}_j|$. For 1-pair: $U_{12} = k\frac{q_1 q_2}{r_{12}}.$ Potential functions: Various cases. Metal: Within a metal region: $\Delta V = V_f - V_i = 0$. Connected metal-regions become an equal potential body.

Spherical shell(Q, R): What is V(r) for r > R and for r < R?

Solid sphere with radius R: For r > R, what is V(r)?

For r < R, $V(r) - V(R) = -\int_{R}^{r} E dr$, with $E = kQ(r/R)^{3}/(4\pi r^{2}) = Cr$. Here C is defined by $kQ/R^{2} = CR$. Long rod (R, Q, linear charge density $\frac{\Delta Q}{\Delta y} = \lambda$). Verify that Gauss law leads to:

for r > R, $E = \frac{\lambda}{2\pi\epsilon_0} \cdot \frac{1}{r}$, and for r < R, $E = \frac{\lambda}{2\pi\epsilon_0} \cdot \left(\frac{r}{R}\right)^2 \cdot \frac{1}{r}$. $V(r) = V(R) - \int_R^r E dr$.

 $\frac{\operatorname{Ring}}{\operatorname{Ring}} (Q, R): V_{ring} = kQ/\rho, \text{ where } \rho = \sqrt{R^2 + z^2}. \text{ Verify that } E_z = -\partial V/\partial z \text{ agrees with Ch16, p639.}$ $\frac{\operatorname{Ring}}{\operatorname{Verify that for large z, it reduces to point charge limit. Inspect the motions of \pm test charges near z=0.}$ $\frac{\operatorname{Parallel plate capacitor}(Q, A, \text{ and gap } s). \text{ Based on Gauss Law, verify that for one plate: } E_1 = \frac{Q/A}{2\epsilon_0}.$ $\frac{\operatorname{Parallel plate capacitor}(Q, A, \text{ and gap } s). \text{ Based on Gauss Law, verify that for one plate: } E_1 = \frac{Q/A}{2\epsilon_0}.$ $\frac{\operatorname{Parallel plate capacitor}(Q, A, \text{ and gap } s). \text{ Based on Gauss Law, verify that for one plate: } E_1 = \frac{Q/A}{2\epsilon_0}.$ $\frac{\operatorname{Parallel plate capacitor}(Q, A, \text{ and gap } s). \text{ Based on Gauss Law, verify that for one plate: } E_1 = \frac{Q/A}{2\epsilon_0}.$ $\frac{\operatorname{Parallel plate capacitor}(Q, A, \text{ and gap } s). \text{ Based on Gauss Law, verify that for one plate: } E_1 = \frac{Q/A}{2\epsilon_0}.$ $\frac{\operatorname{Parallel plate capacitor}(Q, A, \text{ and gap } s). \text{ Based on Gauss Law, verify that for one plate: } E_1 = \frac{Q/A}{2\epsilon_0}.$ $\frac{\operatorname{Parallel plate capacitor}(Q, A, \text{ and gap } s). \text{ Based on Gauss Law, verify that for one plate: } E_1 = \frac{Q/A}{2\epsilon_0}.$ $\frac{\operatorname{Parallel plate capacitor}(Q, A, \text{ and gap } s). \text{ Based on Gauss Law, verify that for one plate: } E_1 = \frac{Q/A}{2\epsilon_0}.$

Magnetic Field. Long wire: The thumb points along the current, direction of **B** pattern is given by RHR1. Horizontal Component of \mathbf{B}_{earth} is ~ 2 × 10⁻⁵T, in much of US.

Biot-Savart law for source $q\mathbf{v}$, $B = \frac{\mu_0}{4\pi} \cdot (q\mathbf{v} \times \hat{\mathbf{r}}) \frac{1}{r^2}$. For a current segment: $\Delta B = \frac{\mu_0}{4\pi} \cdot (I\Delta \mathbf{l} \times \hat{\mathbf{r}}) \frac{1}{r^2}$. In the Drude model, $\Delta Q \bar{\mathbf{v}} = \Delta Q \cdot \frac{\Delta \mathbf{l}}{\Delta t} = \frac{\Delta Q}{\Delta t} \Delta \mathbf{l} = I\Delta \mathbf{l}$, where $\Delta \mathbf{l}$ is the drift distance in time Δt . In terms of electron number density n, $I = q \frac{\Delta N}{\Delta t} = q \frac{nA\bar{\mathbf{v}}\Delta t}{\Delta t} = qi$, where $i = nA\bar{v}$. **Caution**: Among Quest problems, the symbol i is often used for the conventional current. When in doubt, please ask for clarification.

$$\underline{\text{Cross Product}}: \mathbf{A} \times \mathbf{B} = \begin{vmatrix} A_y A_z \\ B_y B_z \end{vmatrix} \hat{i} + \begin{vmatrix} A_z A_x \\ B_z B_x \end{vmatrix} \hat{j} + \begin{vmatrix} A_x A_y \\ B_x B_y \end{vmatrix} \hat{k} = A B \sin \theta \, \hat{\mathbf{n}}. \text{ Direction}: \text{RHR2.}$$

 $\underline{\text{Circular Arc}}: (\mathbf{I}, s = r\theta, \text{ with finite arc length}): \Delta \mathbf{B} \text{ at center, } \Delta \mathbf{B} = \left(\frac{\mu_0}{4\pi}\right) \frac{IR\Delta\theta}{R^2} \hat{\mathbf{n}}, \, \hat{\mathbf{n}} \text{ by RHR2, RHR3.} \\ \underline{\text{Wire segment}}, \text{ Fig18.24 (I, a): Direction: } \Delta \mathbf{B} = -\hat{k}\Delta B_z, \text{ RHR1. } \Delta B_z = \left(\frac{\mu_0}{4\pi}\right) \frac{I\Delta ysin\theta}{r^2}. \ r = \sqrt{a^2 + y^2}. \text{ Use} \\ \underline{\frac{\Delta y}{r^2} = \frac{\Delta\theta}{a}}. \ B_z = \left(\frac{\mu_0}{4\pi}\right) \frac{I}{a}H, \ H = \left(\cos\theta_1 - \cos\theta_2\right) \xrightarrow{symm} 2\cos\theta_1 \xrightarrow{long} 2. \text{ For } \pm \frac{L}{2} \text{ case: } \cos\theta_1 = \frac{L/2}{\sqrt{a^2 + (L/2)^2}}, \text{ p722.} \\ \underline{\text{Ring}} (\mathbf{I}, \mathbf{R}): \ B_z = \left(\frac{\mu_0}{4\pi}\right) \frac{I\cdot 2\pi R}{\rho^2} \cdot \frac{R}{\rho}, \ \rho^2 = R^2 + z^2. \text{ At } z = 0, \ \rho = R. \ B = \left(\frac{\mu_0}{4\pi}\right) \frac{2\pi I}{R}. \\ \text{At } z >> R, \ B = \left(\frac{\mu_0}{4\pi}\right) \frac{2\pi R I}{\rho^3} \equiv \left(\frac{\mu_0}{4\pi}\right) \frac{2\mu}{z^3}, \text{ where dipole moment } \vec{\mu} = A_{loop}I \, \hat{\mathbf{n}}. \ \hat{\mathbf{n}} \text{ is defined by RHR3.} \end{aligned}$

Solenoid: $B = \mu_0 \left(\frac{NI}{2L}\right) (\cos\alpha_1 - \cos\alpha_2) \xrightarrow{long} \mu_0 \frac{NI}{L}$. Clicker 19.3 (skip derivation). B_z vs z curve, Fig. 18.53. Magnetic moment of a bar magnet.

Atomic model: $\mu_{orbit} = \pi R^2 I = \pi R^2 \frac{ev}{2\pi R} = \frac{e}{2m}L$, L = mvR. Ground state of H-atom $L \sim \hbar \sim 10^{-34}$ J s Exp-check: $\mu_{magnet} \sim \mu_{atom} \times (number \ of \ atoms)$. RHS agrees with LHS to within factor of 2 (see p.731).

Ampere's Law: $\oint_{\mathcal{L}} \mathbf{B} \bullet d\mathbf{l} = \mu_0 I_{\mathcal{L}}$. Application: 1) Cyl-symmetry: LHS= $2\pi rB$. Find RHS for long wire, solid wire, cyl-conducting shell. 2) Packed loops(LHSs): Solenoid(Bd), Toroid ($2\pi rB$), current sheet (2Bd).