

1. Consider the quark-antiquark pair production in *Quantum Chromodynamics* (QCD). Specifically, let us focus on the $u\bar{u} \rightarrow d\bar{d}$ process so there is only one tree-level diagram contributing to this process. Draw this diagram and calculate the amplitude, then sum/average the $|\mathcal{M}|^2$ over both spins and *colors* of the final/initial particles and calculate the total cross section. For simplicity, you may neglect the quark masses.

Note that the $u\bar{u} \rightarrow d\bar{d}$ pair production in QCD is very similar to the $e^-e^+ \rightarrow \mu^-\mu^+$ pair production in QED, so the only new aspect of this problem is summing over the colors.

2. Now consider a scalar analogue of QCD or more generally a theory of Yang–Mills fields A_μ^a and complex scalars Φ_i in some representation (r) of the Gauge group G .

(a) Write down the Lagrangian and the Feynman rules of this theory.

Next, consider the annihilation process $\Phi + \Phi^* \rightarrow 2$ gauge bosons. At the tree level, there are four Feynman diagrams contributing to this process.

(b) Draw the diagrams and write down the tree-level annihilation amplitude.

As discussed in class, amplitudes involving the non-abelian gauge fields satisfy a weak form of the Ward identity: *On-shell Amplitudes involving a longitudinally polarized gauge boson vanish, provided all other gauge bosons are transversely polarized.* In other words,

$$\mathcal{M} \equiv e_1^{\mu_1} e_2^{\mu_2} \cdots e_n^{\mu_n} \mathcal{M}_{\mu_1 \mu_2 \cdots \mu_n}(\text{momenta}) = 0$$

when $e_1^\mu \propto k_1^\mu$ but $e_2^\nu k_{2\nu} = \cdots = e_n^\nu k_{n\nu} = 0$.

(c) Verify this identity for the scalar annihilation amplitude.

3. Finally, let us calculate the annihilation cross-section in Scalar QCD. Generally, the tree-level annihilation amplitude (*cf.* the previous problem) can be written in the form

$$\mathcal{M} = F\{T^a, T^b\}_j^i + iG[T^a, T^b]^i_j \quad (1)$$

where j is the ‘color’ index of the scalar particle belonging to some representation (r) of the gauge group G , i is the color index of the scalar anti-particle belonging to the conjugate

representation (\bar{r}), and a and b are the color indices of the gauge bosons belonging to the adjoint representation.

- (a) Show that the annihilation amplitude indeed has form (1) and write down the coefficients F and G as explicit functions of the particles momenta and polarizations.
- (b) Next, let us sum the $|\mathcal{M}|^2$ over the gauge boson's colors and average over the scalar colors. Show that

$$\frac{1}{\dim^2(r)} \sum_{ij} \sum_{ab} |\mathcal{M}|^2 = \frac{C(r)}{\dim(r)} \times (4C(r) |F|^2 + C(\text{adj}) (|G|^2 - |F|^2)). \quad (2)$$

In particular, for scalars in the fundamental representation of the $SU(N)$ gauge group,

$$\frac{1}{N^2} \sum_{ij} \sum_{ab} |\mathcal{M}|^2 = \frac{N^2 - 1}{2N^2} \left(\frac{N^2 - 2}{N} |F|^2 + N |G|^2 \right). \quad (3)$$

- (c) Evaluate F and G in the center of mass frame. Take the gauge boson's polarization vectors $e_{1,2}^\mu$ to be purely spatial and transverse (to the vector particle's momenta).
- (d) Finally, calculate the (polarized, partial) cross-section for the annihilation process.