

1. In a Yukawa theory with $M_s > 2m_f$, the scalar particles are unstable against decay into fermion + antifermion pairs.

Write down the tree-level matrix element $\mathcal{M}(S \rightarrow f + \bar{f})$, sum $|\mathcal{M}|^2$ over final particles' spins, and calculate the total decay rate $S \rightarrow f + \bar{f}$.

2. Consider a Yukawa theory of two Dirac fields $\Psi_1(x)$ and $\Psi_2(x)$ coupled to the same real scalar field $\Phi(x)$:

$$\begin{aligned} \mathcal{L} = & \bar{\Psi}_1(i\not{\partial} - m_1)\Psi_1 + \bar{\Psi}_2(i\not{\partial} - m_2)\Psi_2 + \frac{1}{2}(\partial_\mu\Phi)^2 - \frac{1}{2}M_s^2\Phi^2 \\ & - g_1\Phi\bar{\Psi}_1\Psi_1 - g_2\Phi\bar{\Psi}_2\Psi_2. \end{aligned} \quad (1)$$

At the tree level, calculate the matrix element, the partial cross-section and the total cross-section for elastic scattering of one type of a fermion off the other type, $f_1 + f_2 \rightarrow f_1 + f_2$. Take the initial fermions to be unpolarized and sum over the final fermion's polarizations.

3. Consider a QED-like theory of a massive vector field A^μ coupled to a Dirac field Ψ ,

$$\mathcal{L} = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} + \frac{m_v^2}{2}A^\mu A_\mu + \bar{\Psi}(i\not{\partial} - M_f)\Psi + eA^\mu\bar{\Psi}\gamma_\mu\Psi. \quad (2)$$

The Feynman rules of this theory are similar to those of QED, except for the “photon” propagator: For a massive vector,

$$G_F^{\mu\nu}(x-y) \equiv |0\rangle \mathbf{T}^* \hat{A}^\mu(x) \hat{A}^\nu(y) |0\rangle = -(g^{\mu\nu} + \frac{1}{m_v^2} \partial^\mu \partial^\nu) G_F^{\text{scalar}}(x-y), \quad (3)$$

or in momentum space form,

$$A^\mu \text{---} \text{wavy line} \text{---} A^\nu = \frac{i}{q^2 - m_v^2 + i\epsilon} \left(-g^{\mu\nu} + \frac{q^\mu q^\nu}{m_v^2} \right). \quad (4)$$

Also, the external line factors $e_{\mathbf{k},\lambda}^\mu$ and $e_{\mathbf{k},\lambda}^{*\mu}$ for massive vectors have three allowed polarizations rather than two. (But you don't need them for this exercise.)

- (a) Draw tree diagrams and write down the tree-level amplitude for the elastic scattering of two fermions in this theory. Simplify the amplitude using

$$\bar{u}(p', s') \not{q} u(p, s) = 0 \quad \text{for } q = p' - p. \quad (5)$$

(I shall prove this identity in class this week.)

- (b) Take the non-relativistic limit $|\mathbf{p}| \ll M_f$ for the external fermions and assume the vector field is relatively light, $m_v \ll M_f$. Simplify the scattering amplitude in this limit, compare it to the Born approximation to the non-relativistic scattering of a potential $V(\mathbf{x} - \mathbf{y})$, and show that it agrees with a Yukawa-like potential

$$V(r) = +\frac{e^2}{4\pi r} e^{-m_v r}. \quad (6)$$

Pay attention to the sign of this potential: vector-mediated forces between like charges are repulsive, not attractive.

Now consider the elastic scattering of a fermion and an antifermion.

- (c) Draw the diagrams and write down the exact tree-level amplitude.
- (d) Finally, take the non-relativistic limit as in part (b), simplify the amplitude, and show that it agrees with the Born approximation to scattering off the potential

$$V(r) = -\frac{e^2}{4\pi r} e^{-m_v r}. \quad (7)$$

Again, pay attention to the sign of this potential: The fermion and the antifermion have opposite charges, so they attract rather than repel.