- 1. First of all, finish reading the Witten's paper I assigned last week.
- 2. Consider SQCD with  $N = N_c N_f \ge 2$ . As explained in class, when quark masses are small (or zero) the squarks and antisquarks develop large VEVs and Higgs the  $SU(N_c)$ gauge symmetry down to the SU(N). The low-energy EFT is SYM coupled to the moduli  $\mathcal{M}_{ff'} = \tilde{Q}_f Q'_f$  via moduli-dependence of the SYM's  $\Lambda_I R$ , and the gaugino condensation in the unbroken SYM generates the non-perturbative superpotential for the moduli,

$$W_{n.p.}(\mathcal{M}) = -N \left( \frac{\Lambda_{\text{SQCD}}^{3N_c - N_f}}{\det(\mathcal{M})} \right)^{1/N}, \qquad W_{\text{net}}(\mathcal{M}) = \text{tr}(m\mathcal{M}) + W_{n.p.}(\mathcal{M}).$$
(1)

(a) Suppose all quark flavors have non-zero (but small) masses, so the mass matrix m is invertible. Show that in this case, the SUSY vacuum equations

$$\frac{\partial W_{\text{net}}}{\partial \mathcal{M}_{ff'}} = 0 \quad \forall f, f' \tag{2}$$

have  $N_c$  distinct solutions

$$\mathcal{M}_{ff'} = -(m^{-1})_{ff'} \times \left(\Lambda_{\mathrm{SQCD}}^{3N_c - N_f} \times \det(m)\right)^{1/N_c}, \qquad N_c \text{ branches.}$$
(3)

These solutions correspond to  $N_c$  SUSY vacua of the massive SQCD, which has the same Witten index  $N_c$  as the quark-less SYM.

(b) Now suppose there is only one flavor but it's massless. Your task is to calculate the scalar potential for the modulus  $\Phi = (\mathcal{M} = \tilde{Q}Q)^{1/2}$  and show that for  $\Phi \gg \Lambda_{\text{SQCD}}$  it behaves like a *negative* power of  $|\Phi|$ . This potential does not have any minima at finite  $\Phi$ , hence the theory does not have any stable vacua at all, supersymmetric or otherwise; instead, squark VEVs *run away to infinity*.

Note: Classically, the flat direction  $\Phi$  has  $K(\Phi, \Phi^*) = 2\Phi^*\Phi$  and hence constant metric  $g_{\Phi\bar{\Phi}} = 2$  (prove this!). In the quantum theory, this metric suffers from both perturbative and non-perturbative corrections, but as long as  $\Phi \gg \Lambda$ , the gauge coupling at the Higgs scale is weak and the corrections are small. (c) Finally, consider several massless flavors. For simplicity, assume the moduli matrix is diagonal,  $\mathcal{M} = \text{diag}(\phi_1^2, \phi_2^2, \dots, \phi_{N_f}^2)$ ; classically, this corresponds to diagonal VEVs of the squark and antisquark matrices,  $\langle Q_f^c \rangle = \phi_f \delta_f^c$ ,  $\langle \tilde{Q}_c^f \rangle = \phi_f \delta_c^f$ . Calculate the scalar potential for the  $\phi_f$  (assuming all  $\phi_f \gg \Lambda$ ) and show that they all run away to infinity — there are no stable vacua, supersymmetric or otherwise.