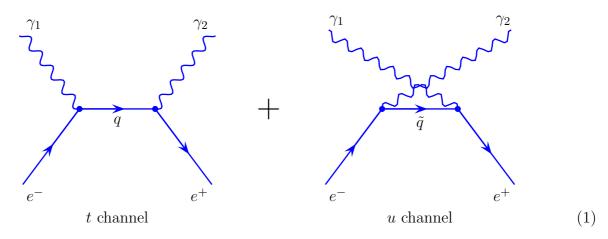
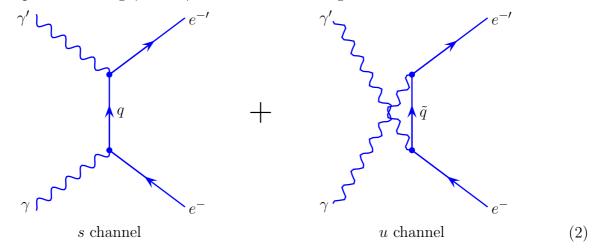
## Annihilation v. Compton Scattering

Annihilation  $e^+e^- \to \gamma\gamma$ : 2 tree diagrams related by Bose symmetry:



Compton scattering  $\gamma e^- \rightarrow \gamma e^-$ : 2 similar tree diagrams



The annihilation and the Compton scattering are related by crossing  $s \leftrightarrow t, i.e.$ 

$$t^a \leftrightarrow s^c, \qquad s^a \leftrightarrow t^c, \qquad u^a \leftrightarrow u^c.$$
 (3)

Hence, the amplitudes of these two processes are analytic continuations of each other as functions of momenta,

$$\mathcal{M}^{\text{annihilation}} = f(s, t, u), \qquad \mathcal{M}^{\text{Compton}} = f(t, s, u)$$
 (4)

for the same analytic function f.

After summing / averaging over fermion spins and photon's polarization, we have

$$\sum_{\lambda_{1},\lambda_{2}} \frac{1}{4} \sum_{s_{-},s_{+}} \left| \mathcal{M}^{\text{annihilation}} \right|^{2} = F(s,t,u),$$

$$\frac{1}{2} \sum_{s,s'} \frac{1}{2} \sum_{\lambda\lambda'} \left| \mathcal{M}^{\text{Compton}} \right|^{2} = -F(t,s,u),$$
(5)

for the same function F. The opposite signs of F in the two formulae follow from crossing an odd number of fermions. These signs are necessary to keep all the spin-summed  $|\mathcal{M}|^2$  positive.

To see where such signs come from, consider some generic process  $X \to Y$  and let

$$f(s,t,u) = \langle Y | \mathcal{M} | X \rangle, \qquad \bar{f}(s,t,u) = \langle X | \mathcal{M} | Y \rangle.$$
 (6)

For real momenta  $\bar{f}$  is the complex conjugate of f, but once we analytically continue to complex momenta, this is no longer true. When we continue to the crossed momenta — which are real but have negative  $p^0$  for the crossed particles — the analytic continuations of the the f and the  $\bar{f}$  functions go back to being complex conjugates of each other up to a sign,

$$\bar{f} = f^* \times (-1)^{\text{\#FC}} \tag{7}$$

where #FC is the number of fermions we have crossed (*i.e.*, made incoming rather than outgoing or vice verse.)