

1. Consider the Seiberg duality below the conformal window: The A theory has  $N_c^A > \frac{2}{3}N_f$  so it is IR-confining rather than IR-conformal, while the B theory has  $N_c^B < \frac{1}{3}N_f$  so it is UV-strong but IR-free. As I explained in class, the B theory emerges as an affective IR theory of the A theory: its “elementary” particles — the  $\Phi^{ff'}$  singlets, the quarks and the antiquarks, and even the gluons — are some kinds of composite particles of the confining A theory. The RG flow of the B theory is not asymptotically free but has a Landau pole at some UV scale  $\Lambda_B$ ; in terms of the A theory, we expect  $\Lambda_B \sim \Lambda_A$ .

Besides the gauge and Yukawa couplings, the B theory may also have non-renormalizable couplings suppressed by the negative powers of the Landau pole, and such couplings may include non-polynomial non-perturbative superpotential terms stemming from some strong interactions at  $E \sim \Lambda_B$ . Thus, the net superpotential of the B theory should look like

$$W_{\text{net}}^B = \lambda \Phi^{ff'} q_f^c \tilde{q}_{cf'} - \mu m_{ff'} \Phi^{ff'} + W_{\text{np}}^B(\Phi, \Lambda_B), \quad (1)$$

where  $m_{ff'}$  are the quark masses of the A theory which translate into the O’Raifeartaigh terms of the B theory.

- (a) Suppose all quarks of the A theory are massive but very light,  $m \ll \Lambda_A$ ; in B theory terms, this translates to  $\mu m \ll \Lambda_B^2$ , so the O’Raifeartaigh terms act as small perturbations. In this regime, the non-perturbative term  $W_{\text{np}}^B(\Phi, \Lambda_B)$  should have all the flavor symmetries of the  $m = 0$  B theory, including the anomalous symmetries which affect the  $\Lambda_B$ .

Show that these symmetries constrain the form of the non-perturbative term to

$$W_{\text{np}}^B(\Phi, \Lambda_B) = \left( \frac{\det(\lambda \Phi)}{(\Lambda^{N_f - 3N_c})_B} \right)^{1/N_c^B} \times \text{a numeric constant.} \quad (2)$$

- (b) Show that for  $\text{rank}(m) = N_f$ , the B theory with superpotential (1) has  $N_f - N_c^B = N_c^A$  supersymmetric vacua. Also show that the  $\langle \mu \Phi^{ff'} \rangle$  VEVs in these vacua agree with the meson VEVs  $\langle \mu \mathcal{M}^{ff'} \rangle$  of the A theory.

2. The rest of this homeworks is primarily a reading assignment:

[Jeffrey A. Harvey's lecture notes on \*Magnetic Monopoles, Duality, and Supersymmetry\*](#).

Lectures 1 and 2 should be fairly easy, so please focus on the harder lectures 3, 4, and 5. Skip lecture 6 for now — this material is not germane to this class. (Although it's rather interesting on its own right, so you might want to read lecture 6 when you have more time — or when your own research encounters the multi-monopole moduli spaces.)

Just for fun, here is a *simple* exercise: Prove the Bogomolny–Prasad–Sommerfeld bound  $M \geq v|Q + i\mu|$  for dyons with both electric and magnetic charges.