- First, a modified textbook problem 9.2(c-e) about the Euclidean path integrals at finite temperature. Questions (a-c) below concern a free scalar field, questions (d-e) concern free fermionic fields, and question (f) is about the free electromagnetic field.
 - (a) Consider a free scalar field in 3+1 dimensions at finite temperature T. Use Euclidean path integral to calculate the partition function and hence the Helmholz free energy. Show that formally

$$\mathcal{F}(T) = \frac{T}{2} \times \operatorname{Tr}\log(-\partial_E^2 + m^2)$$
(1)

where the ∂_E^2 operator acts on functions $(x_1, x_2, x_3, x_4)_E$ which are periodic in the Euclidean time x_4 with period $\beta = 1/T$. Then take the trace in momentum space and the Poisson resummation formula (see homework set#19, eq. (2)) to show that

$$\mathcal{F}(T) = \text{const} + \frac{1}{2} \sum_{\ell = -\infty}^{+\infty} \int \frac{d^4 p_E}{(2\pi)^4} \exp(i\ell\beta p_4) \times \log(p_E^2 + m^2)$$
(2)

$$= \mathcal{F}(0) + \sum_{\ell=1}^{\infty} \int \frac{d^4 p_E}{(2\pi)^4} \exp(i\ell\beta p_4) \times \log(p_E^2 + m^2).$$
(3)

(b) Take the $\int dp_4$ integral in eq. (3). Deform the integration contour so it runs on two 'banks' of a branch cut to show that

$$\int_{-\infty}^{+\infty} \frac{dp_4}{2\pi} \exp(i\ell\beta p_4) \times \log(p_4^2 + E^2) = -\frac{\exp(-\ell\beta E)}{\ell\beta}.$$
 (4)

(c) Now use eqs. (3) and (4) to show that the free energy of a free scalar field above the zero-point energy is

$$\mathcal{F}(T) - \mathcal{F}(0) = \int \frac{d^3 \mathbf{p}}{(2\pi)^3} T \log\left(1 - e^{-\beta E_p}\right) = \int \frac{d^3 \mathbf{p}}{(2\pi)^3} \left(\mathcal{F}_{\text{oscillator}}^{\text{harmonic}}(T, E_p) - \frac{1}{2}E_p\right).$$
(5)

Next, consider a free fermion 0 + 1 dimensions, basically a two-level system in Quantum Mechanics. In the Hamiltonian formulation this means

$$\hat{H} = \omega \hat{\psi}^{\dagger} \hat{\psi}$$
 where $\{\hat{\psi}, \hat{\psi}^{\dagger}\} = 1$ and $\omega = \text{constant} > 0,$ (6)

while in the Lagrangian formulation, $\psi(t)$ and $\psi^*(t)$ are Grassmann-number-valued functions of the time and

$$L_E = \psi^* \times \frac{d\psi}{dt_E} + \omega \times \psi^* \psi.$$
(7)

- (d) Use the path integral to calculate the partition function for both periodic and antiperiodic boundary conditions for the fermionic variables in the Euclidean time, $\psi(t_E + \beta) = \pm \psi(t_E)$. Show that the periodic conditions lead to an unphysical partiction function, while the antiperiodic conditions lead to the correct particition function of a two-level system.
- (e) Now apply the lesson of part (d) to a Dirac fermionic field in 3 + 1 dimensions. Calculate the partition function and hence the free energy using the Euclidean path integral over Dirac fields which are antiperiodic in the Euclidean time, $\Psi(\mathbf{x}, x_4 + \beta) = \Psi(\mathbf{x}, x_4)$.

The last question (f) below involves the path integral for the electromagnetic field. You should answer it only when I have finished explaining how to take such integrals, including the explanation of the Faddeev–Popov determinant and of the gauge-averaging leading to the Feynman gauge condition. If I do not finish this subject the end of the Tuesday 4/11 class, then question (f) is posponed to the next homework.

(f) At finite temperature, the electromagnetic field $A^{\mu}(x)$ — just like any other bosonic field — is periodic in the Euclidean time, $A^{\mu}(\mathbf{x}, x_4 + \beta) = A^{\mu}(\mathbf{x}, x_4)$. Use the path integral to show that formally, the EM free energy is

$$\mathcal{F}(T) = T \times \operatorname{Tr}\log(-\partial_E^2).$$
 (8)

Then recycle calculations from parts (a–c) to show that

$$\mathcal{F}(T) - \mathcal{F}(0) = \int \frac{d^3 \mathbf{p}}{(2\pi)^3} 2T \times \left(1 - e^{-\beta|p|}\right).$$
(9)

The second problem is a big reading assignment: §12.1 of the *Peskin & Schroeder* textbook about integrating out of the short-distance modes and the Wilsonian renormalization. Also, read §12.4 of the Weinberg's book (vol. 1) about the same subjects.

These issues are important, and I wish I could spend a lecture or two explaining them. Alas, the class time is too short, hence this assignment.