Show that in the regime of normal dispersion — *i.e.*, at frequencies not too close to any of the resonances — the group velocity of the EM wave is always less than c. For simplicity, use the low-density approximation

$$\epsilon(\omega) \approx 1 + \frac{Ne^2}{\epsilon_0 m_e} \sum_{i}^{\text{resonances}} \frac{f_i}{\omega_i^2 - \omega^2 - i\omega\gamma_i} \tag{1}$$

as well as $\mu(\omega) \approx 1$.

2. In conducting materials, the EM waves attenuate with distance. For a specific example, consider a uniform material with dielectric constant ϵ , conductivity σ , and negligible magnetism, $\mu = 1$. The attenuating plane wave has general form

$$\mathbf{E}(x, y, z, t) = \vec{\mathcal{E}} \exp(ikz - \kappa z - i\omega t), \quad \mathbf{H}(x, y, z, t) = \vec{\mathcal{H}} \exp(ikz - \kappa z - i\omega t).$$
(2)

(a) Write down formulae for k and κ as functions of ω . Also, relate the electric amplitude $\vec{\mathcal{E}}$ and the magnetic amplitude $\vec{\mathcal{H}}$ to each other.

Now consider a boundary between a conducting material and the vacuum. Suppose an EM wave comes from the vacuum side and hits the boundary head-on.

- (b) Calculate the reflectivity $R = |r|^2$ of the boundary.
- (c) Show that for a good conductor

$$R \approx 1 - \frac{4\pi\delta}{\lambda_0} \tag{3}$$

where λ_0 is the wavelength of the EM wave in the vacuum and δ is the skin-depth of the current of the same frequency in the conductor.

(d) As an example, find the reflectivity of sea water ($\sigma \approx 5 \text{°}/\text{m}$) at an FM radio frequency $\omega = 2\pi \times 100$ MHz.

3. Consider the Goos–Hänchen effect: In a total internal reflection, the reflected ray is displaced sideways relative to the incoming ray as if it's reflected not from the boundary itself but from a small distance behind it.



The key to the Goos–Hänchen effect is the complex reflection coefficient

$$r(\alpha) = \exp(i\phi(\alpha)), \tag{5}$$

its magnitude in a total internal reflection is 1, but the phase depends on the incidence angle α .

(a) Suppose the incident wave has a finite but large width in the direction \perp to the wave within the plane of incidence, for example

$$\mathbf{E}_{i}(x, y, z, t) = \mathcal{E}_{0} \mathbf{e}_{i} \exp\left(ik_{0}(x\sin\alpha + z\cos\alpha) - i\omega t\right) \times \\ \times \exp\left(-\frac{(x\cos\alpha - z\sin\alpha)^{2}}{2a^{2}}\right).$$
(6)

for $a \gg (1/k_0)$. (In my notations, \mathcal{E}_0 is the overall amplitude of the wave and **e** its polarization vector.)

Fourier transform this wave to the \mathbf{k} space, calculate the reflected wave (including its overall phase), then Fourier transform that to the coordinate space. Show that

$$\mathbf{E}_{r}(x, y, z, t) = \mathcal{E}_{0} \mathbf{e}_{r} \exp\left(ik_{0}(x\sin\alpha - z\cos\alpha) - i\omega t\right) \times \int \frac{d\Delta k}{2\pi} A(\Delta k) \times \exp\left(i\Delta k((x\cos\alpha + z\sin\alpha) + i\phi(\mathbf{k}_{0} + \Delta \mathbf{k}))\right)$$
(7)

where $A(\Delta k) = \sqrt{2\pi a} \exp(-a^2 \Delta k^2/2)$.

(b) Perform the Fourier integral in eq. (7) and show that

$$\mathbf{E}_{r}(x, y, z, t) = \mathcal{E}_{0} \mathbf{e}_{r} \exp\left(ik_{0}(x\sin\alpha - z\cos\alpha) - i\omega t\right) \times \\ \times e^{i\phi_{0}} \exp\left(-\frac{(x\cos\alpha + z\sin\alpha - D)^{2}}{2a^{2}}\right)$$
(8)

for the displacement

$$D = -\frac{\partial \phi}{\partial \Delta \mathbf{k}_{\perp}} = -\frac{1}{k_0} \frac{\partial \phi}{\partial \alpha}.$$
 (9)

- (c) Analytically continue the Fresnel equations for the reflection coefficient r to the regime of total internal reflection and calculate its phase ϕ as a function of α . Note two different equations for the in-plane and normal-to-the-plane polarizations of the EM wave.
- (d) Finally, put it all together and calculate the sideways displacement of the reflected wave. Show that

$$D_{\perp} = \frac{2}{k} \frac{\sin \alpha}{\sqrt{\sin^2 \alpha - (n_2/n_1)^4}},$$
 (10)

$$D_{\parallel} = D_{\perp} \times \frac{1}{(1 + (n_1/n_2)^2)\sin^2 \alpha - 1}.$$
 (11)