

1. An inertial frame of reference  $K'$  moves at velocity  $\mathbf{u}$  relative to another inertial frame  $K$ . A particle has velocity  $\mathbf{v}'$  relative to  $K'$ .

- (a) As a warm-up exercise, show that the particle velocity vector relative to the  $K$  frame is

$$\mathbf{v} = \left(1 + \frac{\mathbf{u} \cdot \mathbf{v}'}{c^2}\right)^{-1} \left(\mathbf{u} + \mathbf{v}'_{\parallel} + \frac{1}{\gamma_u} \mathbf{v}'_{\perp}\right) \quad (1)$$

where  $\mathbf{v}'_{\parallel}$  and  $\mathbf{v}'_{\perp}$  are the components of the  $\mathbf{v}'$  vector parallel and perpendicular to the relative velocity  $\mathbf{u}$  of the two frames.

Note: eq. (1) is not symmetric WRT  $\mathbf{v}' \leftrightarrow \mathbf{u}$ , unless the two velocity vectors happen to be parallel to each other.

- (b) Verify that eq. (1) is consistent with the speed of light universality. That is, show that if the particle in question is a photon and  $|\mathbf{v}'| = c$ , then also  $|\mathbf{v}| = c$ , although the directions of the  $\mathbf{v}'$  and  $\mathbf{v}$  vectors are generally different.

Now consider an accelerating particle and the relation between the acceleration vectors  $\mathbf{a}'$  and  $\mathbf{a}$  relative to the frames  $K'$  and  $K$ . Please allow for completely general directions of the vectors  $\mathbf{a}'$ ,  $\mathbf{v}'$ , and  $\mathbf{u}$ .

- (c) Show that

$$\mathbf{a}_{\parallel} = \frac{\left(1 - \frac{\mathbf{u}^2}{c^2}\right)^{3/2}}{\left(1 + \frac{\mathbf{u} \cdot \mathbf{v}'}{c^2}\right)^3} \mathbf{a}'_{\parallel}, \quad (2)$$

$$\mathbf{a}_{\perp} = \frac{\left(1 - \frac{\mathbf{u}^2}{c^2}\right)}{\left(1 + \frac{\mathbf{u} \cdot \mathbf{v}'}{c^2}\right)^3} \left[\mathbf{a}'_{\perp} + \frac{\mathbf{u}}{c^2} \times (\mathbf{a}' \times \mathbf{v}')\right], \quad (3)$$

where  $\mathbf{a}_{\parallel}$  and  $\mathbf{a}_{\perp}$  are the components of the acceleration vector  $\mathbf{a}$  respectively parallel and perpendicular to the relative velocity  $\mathbf{u}$  of the two frames, and ditto for the  $\mathbf{a}'_{\parallel}$  and  $\mathbf{a}'_{\perp}$  components of the  $\mathbf{a}'$  acceleration vector.

2. In the rest frame of a moving conducting medium the current density obeys Ohm's Law  $\mathbf{J}' = \sigma \mathbf{E}'$  where  $\sigma$  is the conductivity and primes denote the rest-frame quantities.

(a) Write the Ohm's Law for a moving conducting medium in a covariant form as

$$J^\mu - \frac{J^\nu U_\nu}{c^2} U^\mu = \frac{\sigma}{c} F^{\mu\nu} U_\nu \quad (4)$$

where  $U^\mu$  is the 4-velocity of the medium,

$$U^\mu = \frac{dX^\mu}{d\tau}, \quad U^0 = \gamma c, \quad U^i = \gamma v^i. \quad (5)$$

(b) Spell out eq. (4) in 3D vector form.

(c) Suppose we are given the electric charge density  $\rho$  in the frame where the medium moves. Show that the electric current density in that frame is (in Gauss units)

$$\mathbf{J} = \rho \mathbf{v} + \gamma \sigma \left[ \mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B} - \frac{\mathbf{v}(\mathbf{v} \cdot \mathbf{E})}{c^2} \right]. \quad (6)$$

(d) Now suppose we know that the moving medium is electrically neutral in its rest frame,  $\rho' = 0$ . Show that in the frame where the medium moves

$$\mathbf{J} = \gamma \sigma \left( \mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B} \right), \quad \rho = \frac{\gamma \sigma}{c^2} (\mathbf{v} \cdot \mathbf{E}). \quad (7)$$

3. Finally, consider a point charge  $Q$  moving at a uniform velocity  $\mathbf{v}$ ; in the worldline terms,

$$x_{\text{charge}}^\mu(\tau) = U^\mu \tau. \quad (8)$$

Let's measure the electric and magnetic fields of this charge at some spacetime point  $x^\mu$ .

- (a) Use the Lorentz covariant of the EM field tensor  $F^{\mu\nu}$  — and only its covariance but no specific EM equations — and the translational invariance of space and time to argue that

$$F^{\mu\nu}(x) = Q(X^\mu U^\nu - X^\nu U^\mu) \times f((X \cdot U), (X \cdot X)) \quad (9)$$

where  $X^\mu = x^\mu - x_{\text{charge}}^\mu$  and  $f$  is some function of the Lorentz scalars  $(X \cdot U) = X^\lambda U_\lambda$  and  $(X \cdot X) = X^\lambda X_\lambda$ .

Note: there is another Lorentz scalar at play, namely  $(U \cdot U) = U^\lambda U_\lambda$ , but its value is a constant  $(U \cdot U) = c^2$ ,

The specific formula for the EM field tensor is

$$F^{\mu\nu}(x) = \frac{Q}{c} \times \frac{(X^\mu U^\nu - X^\nu U^\mu)}{\left(\frac{1}{c^2}(X \cdot U)^2 - (X \cdot X)\right)^{3/2}}. \quad (10)$$

In this formula  $X^\mu = x^\mu - x_{\text{charge}}^\mu(\tau)$ , but the  $\tau$  of the charge can affect the EM field tensor at any given measurement point  $x^\mu$ . So for simplicity we may set  $\tau = 0$  and rewrite eq. (10) as

$$F^{\mu\nu}(x) = \frac{Q}{c} \times \frac{(x^\mu U^\nu - x^\nu U^\mu)}{\left(\frac{1}{c^2}(x \cdot U)^2 - (x \cdot x)\right)^{3/2}}. \quad (11)$$

- (b) Verify this statement.
- (c) Show that for a charge at rest, eq. (10) yields the Coulomb electric field and zero magnetic field.
- (d) For a moving charge, spell out eq. (10) for the electric and magnetic fields in 3-vector notations.