

E3/1

2 boson Hilbert space

$$\Psi(\vec{x}_1, \vec{x}_2) = \Psi(\vec{x}_2, \vec{x}_1)$$

$N$  boson wave functions

$$\Psi(\vec{x}_1, \dots, \vec{x}_N) = \Psi(\text{any permutation}) \\ \text{of } \vec{x}_1, \dots, \vec{x}_N$$

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$$\hat{\Phi}(\vec{x}) = \sum_{\vec{p}} L^{-3/2} e^{-i\vec{p} \cdot \vec{x}/\hbar} \hat{a}_{\vec{p}}$$

$$\hat{\Psi}^{\dagger}(\vec{x}) = \sum_{\vec{p}} L^{-3/2} e^{+i\vec{p} \cdot \vec{x}/\hbar} \hat{a}_{\vec{p}}^{\dagger}$$

$\hat{\Psi}^{\dagger}$  creates an extra atom  
@ point  $\vec{x}$

$\hat{\Psi}$  destroys an atom @  $\vec{x}$ .

# Field-Particle Duality

Exactly same Quantum theory  
but 2 different classical  
(or semi-classical) limits

- $\alpha$ ) classical Fields
- $\beta$ ) classical particles.

QED     $\alpha$ ) classical EM fields  
          $\beta$ ) Photons.

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$\alpha$ ) classical atoms

↓  
Quantum theory

↓  
 $\beta$ ) classical  $\phi(x)$  &  $\phi^*(x)$  fields.

→ non zero  $\phi(x)$  in a

Bose-Einstein condensate.

$\phi(x)$  describe the hydrodynamics  
of the superfluid component.

Bosonic atoms like  $^4\text{He}$  @ zero temperature

Naturally: all atoms in same state  $|\vec{p}=0\rangle$

$$|\text{naive}\rangle = |n_0=N, \text{ other } n_{\vec{p}}=0\rangle$$

Trouble: long-distance quantum correlations

$$\begin{aligned} G(x,y) &= \langle \text{naive} | \hat{\Psi}^\dagger(x) \hat{\Psi}(y) | \text{naive} \rangle \\ &= \langle \text{naive} | \hat{\Psi}^\dagger(x) | \text{naive} \rangle \langle \text{naive} | \hat{\Psi}(y) | \text{naive} \rangle \\ &= \frac{N}{\text{volume}} \neq 0 \text{ even for } |x-y| \rightarrow \infty \end{aligned}$$

→ bad.

Simplest solution: coherent state

$$|\text{coherent}\rangle = e^{-\bar{n}/2} \exp(\sqrt{\bar{n}} \hat{a}_0^\dagger) |vac\rangle.$$

$$\langle \hat{N} \rangle = \bar{n}$$

$$\Delta N = \sqrt{\bar{n}} \ll \langle \hat{N} \rangle$$

$$\hat{\Psi}(x) |\text{coherent}\rangle = \sqrt{\bar{n}} |\text{coherent}\rangle$$

$$\bar{n} = \frac{\bar{N}}{\text{volume}}$$

# Landau - Ginzburg Model

$$V_2(\vec{x}-\vec{y}) = \lambda \delta^{(3)}(\vec{x}-\vec{y})$$

good for BEC of heavy atoms

starting approx for  ${}^4\text{He}$

$$\hat{H} - \mu \hat{N} = \int d^3x \left( \frac{\hbar^2}{2m} \nabla \psi^\dagger \cdot \nabla \psi + \frac{\lambda}{2} \psi^\dagger \psi^\dagger \psi \psi - \mu \psi^\dagger \psi \right)$$

coherent states  $\rightarrow$  classical fields

$$\phi = \langle \hat{\psi} \rangle \text{ and } \phi^* = \langle \hat{\psi}^\dagger \rangle$$

$$\phi(x) = \langle \hat{\psi}(x) \rangle, \quad \phi^*(x) = \langle \hat{\psi}^\dagger(x) \rangle$$

$$\text{Free Energy } F = E - \mu N = \int d^3x \left( \frac{\hbar^2}{2m} |\nabla \phi|^2 + \frac{\lambda}{2} |\phi|^4 - \mu |\phi|^2 \right)$$

@ Temperature = 0, minimize F.

$$\hookrightarrow \nabla \phi(x) \equiv 0 \rightarrow \phi(x) = \text{const.}$$

$$\hookrightarrow \text{minimize } \frac{\lambda}{2} |\phi|^4 - \mu |\phi|^2 = \frac{\lambda}{2} \left( |\phi|^2 - \frac{\mu}{\lambda} \right)^2 - \frac{\mu^2}{2\lambda}$$

For  $\mu < 0$ , minimum @  $\phi(x) \equiv 0$   
vacuum

For  $\mu > 0$ , minimum @  $|\phi(x)|^2 = \frac{\mu}{\lambda}$

$$\rightarrow \phi(x) \equiv \sqrt{\frac{\mu}{\lambda}} e^{i \text{phase}}$$

any const. phase.

undergraduate stat mech

→ free atoms,  $\lambda = 0$

↳ BEC requires  $\mu < 0$  @  $T > 0$

$$\mu = 0 \text{ @ } T = 0.$$

real life: interaction  $\lambda > 0$  requires

$$\mu > 0 \text{ @ } T = 0$$

$$\mu = \lambda \bar{n} \text{ density}$$

Moving BEC for  $\phi(x) \neq \text{const.}$

$$\hat{N} = \int d^3x \hat{n} \quad \hat{n}(x) = \hat{\psi}^\dagger(x) \hat{\psi}(x)$$

→ classical  $n(x) = |\phi(x)|^2$

is the BEC density.

[in liq He,  $|\phi|^2 = \text{superfluid density}$ ]

$$\vec{P}_{\text{net}} = \int d^3x \hat{\psi}^\dagger(-i\hbar) \vec{\nabla} \hat{\psi}$$

$$\rightarrow \int d^3x \phi^* (-i\hbar) \nabla \phi = -\frac{i\hbar}{2} \int d^3x (\phi^* \nabla \phi - \phi \nabla \phi^*)$$

$$= \frac{\hbar}{2m} \int d^3x \text{Im}(\phi^* \nabla \phi)$$

→ momentum density

$$\frac{\vec{P}}{\text{vol}} = \hbar \text{Im}(\phi^* \nabla \phi) = \hbar |\phi|^2 \vec{\nabla} \text{phase}(\phi(x))$$

$$\frac{\vec{p}}{vol.} = \frac{\hbar}{i} |\phi|^2 \nabla \text{phase}(\phi(\vec{x}))$$

$$\frac{\vec{p}}{vol.} = \rho \vec{v} = m n \vec{v}$$

$\vec{v}$ : flow velocity

$$m n \vec{v} = \frac{\hbar}{i} \cdot (|\phi|^2 = n) \cdot \nabla \text{phase}(\phi)$$

$$\vec{v} = \frac{\hbar}{m} \nabla \text{phase}(\phi)$$

~~$|\phi|^2$~~   $|\phi(\vec{x})|^2 \rightarrow$  superfluid density

$\nabla \text{phase}(\phi) \rightarrow \frac{m}{\hbar} \cdot \text{velocity}$