

1. Let's start with the spin states of a silver atom. As discussed in class, silver atoms have only two independent spin states, so in this problem we are dealing with a two-dimensional Hilbert space. However, measuring the atom's magnetic moment  $\mathbf{m}$  in different directions gives us different bases for this Hilbert space:  $(|Z+\rangle, |Z-\rangle)$  are states of definite  $m_z = \pm m_B$  (where  $m_B = e\hbar/2M_e c$  is the Bohr magneton),  $(|X+\rangle, |X-\rangle)$  are states of definite  $m_x = \pm m_B$ ,  $(|Y+\rangle, |Y-\rangle)$  are states of definite  $m_y = \pm m_B$ , and likewise for the more general directions of the magnetic moment.

Interpreting the Stern-Gerlach experiments with multiple magnetic gaps in terms of the Born rule of quantum probabilities tells us that

$$\begin{aligned} \langle X+ | X- \rangle &= \langle Y+ | Y- \rangle = \langle Z+ | Z- \rangle = 0, \\ |\langle Z \pm | X \pm \rangle|^2 &= |\langle Z \pm | Y \pm \rangle|^2 = |\langle X \pm | Y \pm \rangle|^2 = \frac{1}{2}. \end{aligned} \quad (1)$$

- (a) Use eqs. (1) to show that after a physically-irrelevant change of the overall phases of the ket-vectors  $|Z\pm\rangle$ ,  $|X\pm\rangle$ , and  $|Y\pm\rangle$ , the six quantum states become related to each other as

$$|X\pm\rangle = \sqrt{\frac{1}{2}} |Z+\rangle \pm \sqrt{\frac{1}{2}} |Z-\rangle, \quad |Y\pm\rangle = \sqrt{\frac{1}{2}} |Z+\rangle \pm i\sqrt{\frac{1}{2}} |Z-\rangle. \quad (2)$$

- (b) Construct the operators  $\hat{m}_x$ ,  $\hat{m}_y$ , and  $\hat{m}_z$  for the three components of the atom's magnetic moment and use eqs. (2) to write down the matrices of those operators in the  $(|Z+\rangle, |Z-\rangle)$  basis.
- (c) For an arbitrary direction pointed by a unit vector  $\mathbf{n}$ , the  $n^{\text{th}}$ -component of the atom's magnetic moment is  $m_n = \mathbf{n} \cdot \mathbf{m}$ . Write down the matrix of the operator  $\hat{m}_n = n_x \hat{m}_x + n_y \hat{m}_y + n_z \hat{m}_z$  in the  $(|Z+\rangle, |Z-\rangle)$  basis. Calculate the eigenvalues of this matrix and explain the physical meaning of your result.

2. Brush up your knowledge of basic complex analysis. Focus on analytic functions  $f(z)$  and contour integrals  $\int_{\Gamma} f(z) dz$  over generic contours  $\Gamma$  in the complex plane; make sure you understand which deformations of the contour  $\Gamma$  leave the integral invariant and which do not.

If you have never studied complex analysis before, a good introductory textbook is *Complex Variables* in the *Schaum Outlines* series, by Spiegel, Lipschutz, Schiller, and Spellman; the PMA library has a few copies. For the purposes of this homework focus on chapter 4; then read the rest of the book when you have time.

As a simple test of your knowledge, make sure you understand why for any complex  $\alpha$  with  $\text{Re } \alpha > 0$  and any complex  $\beta$ ,

$$\int_{\substack{\text{real} \\ \text{axis}}} \exp(-\alpha(z - \beta)^2) dz = \sqrt{\frac{\pi}{\alpha}} \quad (3)$$

where the integration contour is the whole real axis from  $-\infty$  to  $+\infty$ .

3. Consider a one-dimensional quantum particle with a Gaussian wave function

$$\Psi(x) = C e^{ax^2 + bx} \quad (4)$$

where  $a$ ,  $b$  and  $C$  are some complex numbers. Note that the discussion in class was limited to the case of real  $a < 0$  but in this exercise we allow for any complex  $a$  with  $\text{Re } a < 0$ .

- (a) Calculate the norm  $\int dx |\Psi(x)|^2$  of this wave function.
- (b) Calculate the momentum-space wave function  $\tilde{\Psi}(p)$  and show that it also has a Gaussian form

$$\tilde{\Psi}(p) = \tilde{C} e^{\tilde{a}p^2 + \tilde{b}p} \quad (5)$$

for some parameters  $\tilde{a}$ ,  $\tilde{b}$ , and  $\tilde{C}$ . Write down these parameters as explicit functions of  $a$ ,  $b$ , and  $C$ .

- (c) Verify that the momentum-space and the coordinate-space wave functions have the same norm.

- (d) Calculate the expectation values  $\langle x \rangle$ ,  $\langle p \rangle$  and the uncertainties  $\Delta x$  and  $\Delta p$ .
- (e) Show that for any Gaussian wave function  $\Delta X \cdot \Delta P \geq \hbar/2$  and that the equality is achieved whenever  $a$  is real.
- (f) The Hamiltonian of a free non-relativistic particle is

$$\hat{H} = \frac{\hat{p}^2}{2M}, \quad (6)$$

so in the momentum space

$$\hat{H}\tilde{\Psi}(p) = \frac{p^2}{2M} \times \tilde{\Psi}(p). \quad (7)$$

Solve the time-dependent Schrödinger equation in the momentum space, and show that if the momentum-space wave function has form (5) at time  $t = 0$ , then it also has form (5) at any time  $t$ , albeit for some different parameters  $\tilde{a}(t)$ ,  $\tilde{b}(t)$ , and  $\tilde{C}(t)$ .

Note that this immediately implies that if the coordinate-space wave function has form (4) at time  $t = 0$ , then it has form (4) at any time  $t$ , albeit for time-dependent parameters  $a(t)$ ,  $b(t)$ , and  $C(t)$ .

- (g) Evaluate the time-dependence of the expectation values  $\langle x \rangle$ ,  $\langle p \rangle$  and the uncertainties  $\Delta x$  and  $\Delta p$  and explain the physical meaning of your results.  
For simplicity, assume real  $a_0 < 0$  at the initial time  $t = 0$ .