

E5/1

Complex scalar field $\phi(x)$

U(1) gauge symmetry

$$\phi(x) \rightarrow e^{i\theta(x)} \phi(x), \text{ under } \theta(x) \text{ @ each } x$$

Lattice, we sample $\phi(x)$ @ $x^\mu = a n^\mu$

$$\text{Gauge: } \phi(\vec{n}) \rightarrow e^{i\theta(\vec{n})} \phi(\vec{n})$$

under $\theta(\vec{n})$ @ each lattice site $\vec{n}$

Continuum: covariant $\mathcal{D}_\mu \phi(x) = \partial_\mu \phi(x) + i A_\mu(x) \phi(x)$

$$A'_\mu(x) = A_\mu(x) - \partial_\mu \theta(x)$$

$$(\mathcal{D}_\mu \phi)'(x) = e^{i\theta(x)} \mathcal{D}_\mu \phi(x)$$

$$\text{Lattice: } \partial_\mu \phi(x) \rightarrow \frac{1}{a} (\phi(\vec{n} + \hat{\mu}) - \phi(\vec{n}))$$

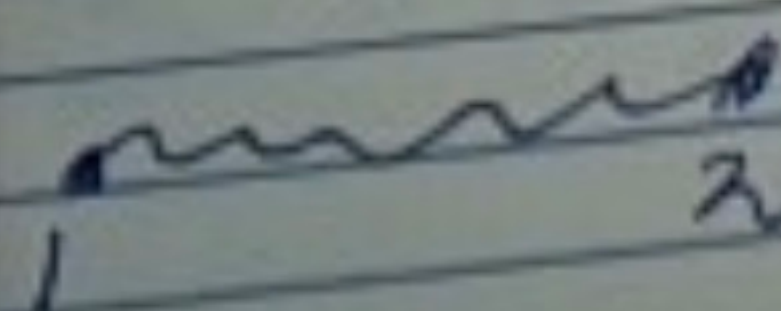
diff phase transformation

to make covariant derivative,

need a parallel transport from

$\vec{n}$ to $\vec{n} + \hat{\mu}$

Take any path for the continuum



$$\text{let } U(\text{path}) = \exp\left[-i \int_1^2 A_\mu dx^\mu\right]$$

along the path

E5/2

Gauge trans:

$$\int_1^2 A_\mu dx^\mu \rightarrow \text{itself} - \int_1^2 \partial_\mu \theta dx^\mu \\ = \text{itself} - \theta(2) + \theta(1).$$

$$W(\text{path}) \rightarrow \text{itself} \times e^{i\theta(2)} e^{-i\theta(1)}.$$

$$\& W(1 \rightarrow 2) \phi(1) \rightarrow e^{i\theta(2)} W(1 \rightarrow 2) \phi(1).$$

$\hookrightarrow$ transforms like $\phi(2)$.

$$\phi(2) = W(1 \rightarrow 2) \phi(1) \hookrightarrow \text{covariant.}$$

Lattice: $\forall$ lattice link from $\vec{u}$ to $\vec{u} + \vec{\mu}$

let $W_\mu(\vec{u})$ be $W(\text{that link})$

$$= \exp(-i \int_{\vec{u}}^{\vec{u} + \vec{\mu}} A_\mu dx^\mu).$$

$$= 1 - i a A_\mu(\vec{u}) + O(a^2).$$

$$\text{Lattice Dirac} = \frac{1}{a} [\phi(\vec{u} + \vec{\mu}) - W_\mu(\vec{u}) \phi(\vec{u})]$$

Gauge-invariant lattice action

for ϕ

$$S = a^4 \sum_{\vec{u}} \left\{ \frac{1}{a^3} \sum_{\vec{\mu}} \left| \phi(\vec{u} + \vec{\mu}) - W_\mu(\vec{u}) \phi(\vec{u}) \right|^2 \right. \\ \left. - m^2 |\phi(\vec{u})|^2 - \frac{1}{4} |F_{\mu\nu}|^2 \right\}$$

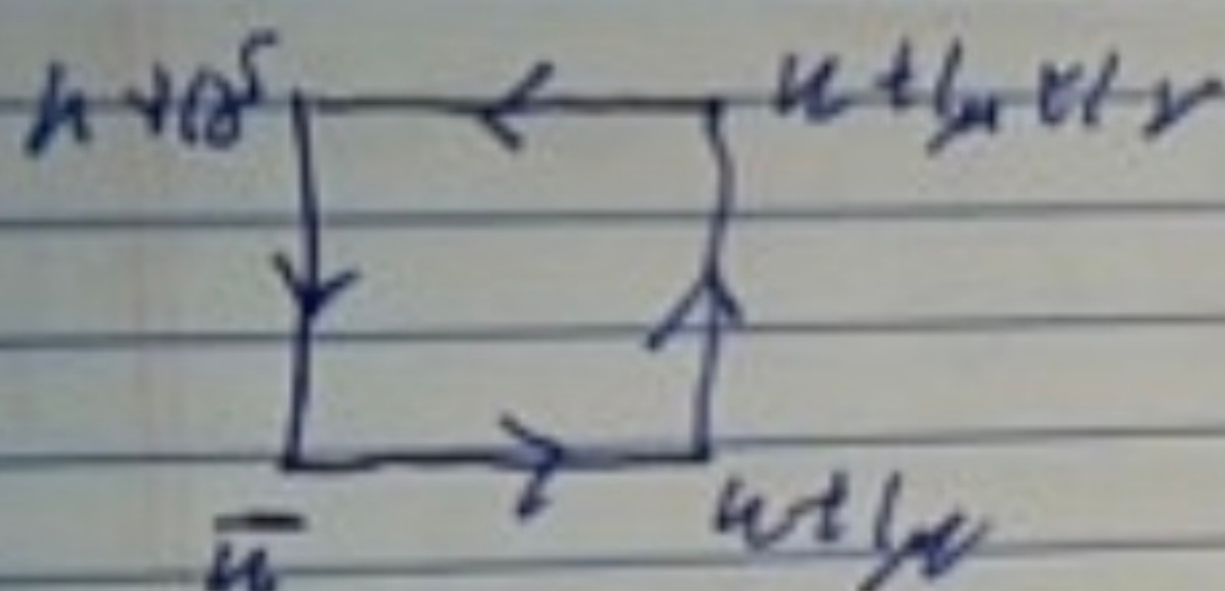
E5/3

Gauge field $A_\mu(x)$ is sampled
via link variables $w_\mu(\vec{u})$.

$$|w_\mu(\vec{u})| = 1.$$

$$\text{Gauge } w_\mu(\vec{u}) \rightarrow e^{i\theta(\vec{u}+\vec{\mu})} w_\mu(\vec{u}) e^{-i\theta(\vec{u})}$$

Action: Take a unit plaquette



Define $w_\mu(\vec{u})$

$$= w_{21}^{-1}(\vec{u}) w_{12}^{-1}(\vec{u}+\vec{\mu}) w_{23}(\vec{u}+\vec{\mu}) w_{32}(\vec{u})$$

w_μ gauge invariant.

$$w_{\mu\nu} = w_{\nu\mu}^{-1}$$

in terms of continuous $A_\mu(x)$.

$$w_\mu(\vec{u}) = \oint_{\text{link}} \prod \exp(-i \int A_\mu dx^\mu)$$

$$= \exp(-i \oint A_\mu dx^\mu)$$

$$= \exp(-i \int \int F_{\mu\nu} dx^\mu dx^\nu)$$

$$a \rightarrow 0 \text{ limit } \oint = a^2 F_{\mu\nu}(\vec{u}) + O(a^3)$$

$$w_\mu(\vec{u}) \approx \exp(-ia^2 F_{\mu\nu}(\vec{u}))$$

E5/4

$$2 - w_{\mu\nu} - w_{\mu\nu}^{-1} \approx 2 - 2 \cos(a^2 \nabla_\mu \nabla_\nu) \\ = + a^4 (\nabla_\mu \nabla_\nu)^2 \quad (\text{no } \sum_\mu)$$

Lattice action

$$S = \frac{1}{4a^2} \sum_\mu \sum_{\mu \neq \nu} \approx \frac{a^4}{4a^2} \sum_\mu \sum_{\mu \neq \nu} (\nabla_\mu \nabla_\nu)^2$$

$$\rightarrow \frac{1}{4a^2} \int d^4x \nabla_\mu \nabla_\nu \nabla_\mu \nabla_\nu = S_E(A_\mu)$$

$$S_E^{\text{lattice}} = \frac{1}{4a^2} \sum_\mu \sum_{\mu \neq \nu} (2 - w_{\mu\nu} - w_{\mu\nu}^{-1})$$

$\hookrightarrow$ Eucl EM action in terms of the
link variables $w_\mu(\vec{u})$

Fund. integral

$$\int \mathcal{D}[\phi(x)] \int \mathcal{D}[\phi^*(x)] e^{-S_E}$$

$$\rightarrow \text{cont limit of } (a^4) \times \prod_{\vec{u}} \int d\phi(\vec{u}) d\phi^*(\vec{u}) e^{-S_E^{\text{cont.}}}$$

lattice variables: complex $\phi(\vec{u})$ at each
site.

For A_μ , lattice variables are

the $w_\mu(\vec{u})$ constrained to $|w_\mu(\vec{u})| = 1$

or Higgs phase: $w_\mu(\vec{u}) \rightarrow e^{i\theta_\mu(\vec{u})}$ where $w_\mu(\vec{u}) = e^{i\theta_\mu(\vec{u})}$

ES15]

Lattice funct. integral

$\oint \mathcal{D}(\psi, \bar{\psi}) \rightarrow$ cont. limit of

$\prod_{\text{sites } n \in V} \oint d\text{measure}(\psi_n, \bar{\psi}_n)$
 each $\psi_n, \bar{\psi}_n$ spend a
 cert. amt in $\mathbb{C}$

measure $d(\psi_n, \bar{\psi}_n)$ must be gauge inv.

$$d\text{measure}(\psi_n, \bar{\psi}_n) = \frac{1}{2\pi} d\text{phase}(\psi_n, \bar{\psi}_n)$$

$$\Phi_{\text{Gauss}}: d\text{phase}(\psi_n, \bar{\psi}_n)$$

$$\rightarrow d(\text{phase}(\psi_n, \bar{\psi}_n) + \theta(N+1) - \theta(N))$$

$$= d\text{phase}(\psi_n, \bar{\psi}_n): \text{gauge inv.}$$

$\frac{1}{2\pi}$ prefactor makes gauge redundancy
 harmless.

$$\prod_{\text{all sites } n} \oint d\text{phase}(\psi_n, \bar{\psi}_n) \Rightarrow \prod_{\text{indep links}} \oint \frac{1}{2\pi} d\text{phase}(\psi_n, \bar{\psi}_n)$$

$$\times \prod_{\text{gauge redundant links}} \oint \frac{1}{2\pi} d\text{phase}(\psi_n, \bar{\psi}_n) \quad \left. \begin{array}{l} \text{is harmless} \\ 1. \end{array} \right\}$$

$$\oint d\text{phase} = 1$$

ES/6

on the lattice, we don't need
to fix a gauge!

Non Abelian gauge Theories.

Let $G = SU(N)$, $\phi^c \in \text{fundamental}$.

$$\Phi(x) = \begin{pmatrix} \phi_1(x) \\ \vdots \\ \phi_N(x) \end{pmatrix}$$

$$V(\phi^\dagger \phi) = v_0^2 \phi^\dagger \phi + \frac{\lambda}{4} (\phi^\dagger \phi)^2$$

$\hookrightarrow$ scalar invariant.

under local $SU(N)$,

$$\phi^c(x) \rightarrow U^c_j(x) \phi^c(x)$$

$U^c_j(x)$ is an $SU(N)$ matrix

indep. matrix $\forall x$.

Lattice: @ each lattice site

$$\Phi(\vec{r}) = \begin{pmatrix} \phi_1(\vec{r}) \\ \vdots \\ \phi_N(\vec{r}) \end{pmatrix}$$

Gauge tr. $\phi^c(\vec{r}) \rightarrow U(\vec{r}) \phi^c(\vec{r})$

indep. scalar matrix $U(\vec{r})$

@ each $\vec{r}$.

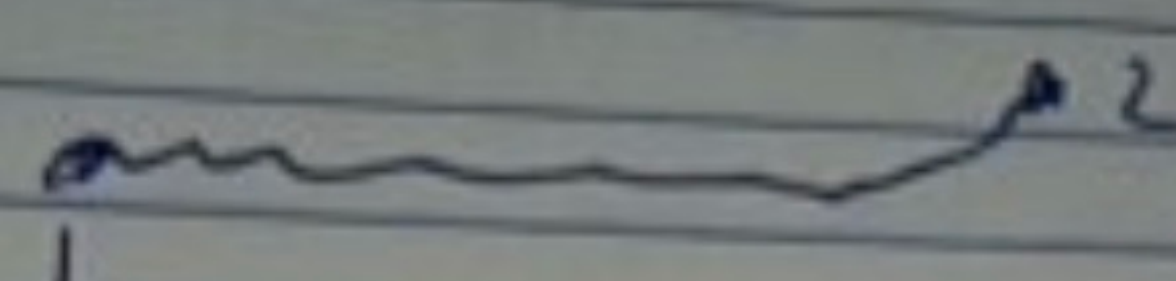
② Covariant derivative

$$D_\mu \phi^c(x) = \partial_\mu \phi^c(x) + i A_\mu^c(x) \phi^c(x)$$

A_μ : matrix-valued gauge field.

ES/7]

$$A_0(x) = U(x) A_m(x) U^\dagger(x) + i \partial_\mu U(x) U^\dagger(x).$$

Take a path 

$$\text{define } W(\text{path}) = P\text{-exp} \left(-i \int_{\text{path}} A_\mu dx^\mu \right).$$

↓
path-ordered exp.

path $x^\mu(t)$ $t: \text{time}$

$$x(0) = x_1, x(1) = x_2.$$

$$W = 1 + \sum_{K=1}^{\infty} (-i)^K \int \int \dots \int_{0 \leq t_1 < t_2 < \dots < t_K \leq 1} dt_1 \dots dt_K A_\mu(x_1) A_\mu(x_2) \dots A_\mu(x_K).$$

Theorem 1: For any kernel, kernels $A_\mu(x)$ $W(\text{path})$ is an $SU(N)$ matrix

Theorem 2: under gauge transformations

$$W(\text{path } 1 \rightarrow 2) \rightarrow U(2) W(\text{path}) U^\dagger(1).$$

$$\hookrightarrow W(1 \rightarrow 2) \Phi(1) \text{ has form like } \Phi(2).$$

E5/8]

On the lattice, we sample $A_\mu(x)$
via $W_\mu(\vec{x}) = W(\text{path from } \vec{x} \text{ to } \vec{x} + \hat{\mu})$
all $W_\mu(\vec{x})$ are $U(1)$ matrices.

Covariant $D_\mu \phi(x) \rightarrow$

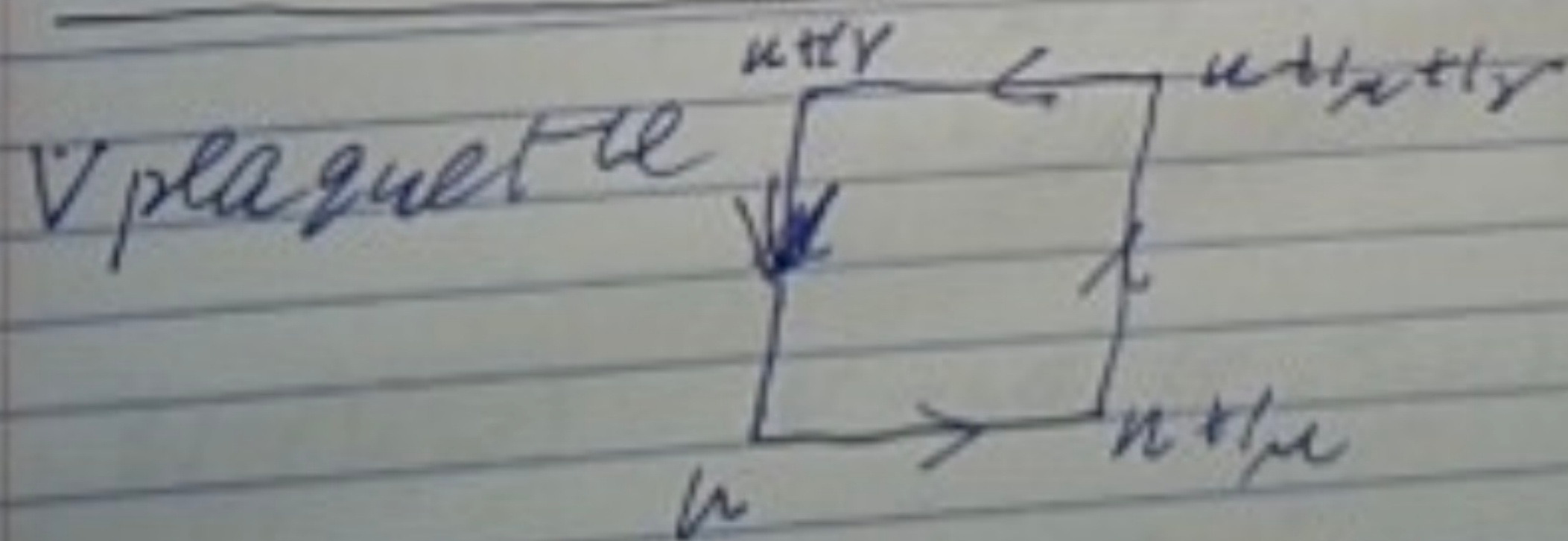
$$\frac{1}{a} [\phi(\vec{x} + \hat{\mu}) - W_\mu(\vec{x}) \phi(\vec{x})]$$

$$S^{\text{Lattice}}(\phi) = a^4 \sum_{\vec{x}} \left\{ V(\phi(\vec{x})^\dagger \phi(\vec{x})) \right. \\ \left. + \frac{1}{2} \sum_{\vec{\mu}} (\phi_{\vec{x}+\hat{\mu}}^\dagger - \phi_{\vec{x}}^\dagger W_\mu(\vec{x})) (\phi_{\vec{x}} - W_\mu^\dagger(\vec{x}) \phi_{\vec{x}+\hat{\mu}}) \right\}$$

Gauge invariant.

Lattice action for gauge fields

$\rightarrow$ want $\frac{1}{2g^2} \text{Tr}(F_\mu F_\mu)$ in terms
of $W_\mu(\vec{x})$.



$$W_{\mu\nu}(\vec{x}) = W_\nu^{-1}(\vec{x}) W_\mu(\vec{x}) W_\nu(\vec{x} + \hat{\mu}) W_\mu^{-1}(\vec{x} + \hat{\nu})$$

$\equiv$
 $W(\square)$ in this order.

25/9

Gauge transformation

$$W_{\mu\nu}(\vec{u}) \rightarrow U(\vec{u}) W_{\mu\nu}(\vec{u}) U^\dagger(\vec{u})$$

$\text{tr}(W_{\mu\nu}(\vec{u}))$ is gauge invariant.

Last, let's:

$$W_{\mu\nu}(\vec{u}) = P\text{-exp}(-i \oint A_\mu dx^\mu)$$

$$= 1 - i a^2 F_{\mu\nu}(\vec{u}) + O(a^3)$$

$$\left[F_{\mu\nu}^2 \partial_\mu A_\nu - \partial_\mu A_\nu + i [A_\mu, A_\nu] \right]$$

$$W_{\mu\nu}(\vec{u}) \approx 1 - i a A_\mu(\vec{u})$$

$$W_{\mu\nu}(\vec{u}) = (1 + i a A_\mu(\vec{u})) (1 + i a A_\nu(\vec{u} + \vec{a}))$$

$$(1 - i a A_\mu(\vec{u} + \vec{a})) (1 - i a A_\nu(\vec{u}))$$

$$= 1 - i a^2 \partial_\mu A_\nu + i a^2 \partial_\nu A_\mu$$

$$= a^2 [A_\mu, A_\nu] + O(a^3)$$

$$\text{tr}(W_{\mu\nu}) = N - i a^2 \text{tr}(F_{\mu\nu})$$

$$+ \frac{a^4}{2} \text{tr}(F_{\mu\nu}^2)$$

$$\text{tr}(W_{\mu\nu}) \rightarrow \text{tr}(W_{\mu\nu}^{-1})$$

$$= a^4 \text{tr}(F_{\mu\nu}^2)$$

E5/10

$$S^{\text{lathe}}(\phi_n) = \frac{1}{2g^2} \sum_n \sum_{\mu, \nu} (\phi_n - \text{tr}(\phi_{\mu\nu} + \phi_{\mu\nu}^{-1}))$$

→ cont. limit $\frac{1}{2g^2} \int d^4x \text{tr}(F_{\mu\nu} F_{\mu\nu})$

S^{lathe} : gauge invariant.

Functional integral

$\int \mathcal{D}(\phi(x)) e^{-S_E}$ is cont. limit of

$$\prod_{\text{links } \ell \in G} \int_{\text{lathe}} d\text{measure}(u_\ell) e^{-S_E}$$

& since $\int_{\text{lathe}} d\text{measure}(u_\ell)$ spans the group manifold of $S(U)$.

$d\text{measure}(u_\ell)$ must be invariant wrt

$$u \rightarrow u' = U(u)uU^{-1}(u).$$

Fortunately, & compact simple group T unique invariant measure.

$$d(u) = d(U_1 u U_2^{-1})$$

for any fixed U_1, U_2 .

ES/II For $U(2)$, $SU(2)$

group manifold is a 3-sphere S^3

$$U = \alpha + i \vec{\beta} \cdot \vec{\sigma}$$

$\vec{\sigma} \rightarrow$ Pauli

$$\alpha^2 + \vec{\beta}^2 = 1$$

$\rightarrow (\alpha, \vec{\beta})$ spans S^3

inv. measure is the ³volume measure
on S^3 .